

Synthetic gauge field and synthetic dimension in interacting ultracold Fermi gases

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Work in collaboration with...



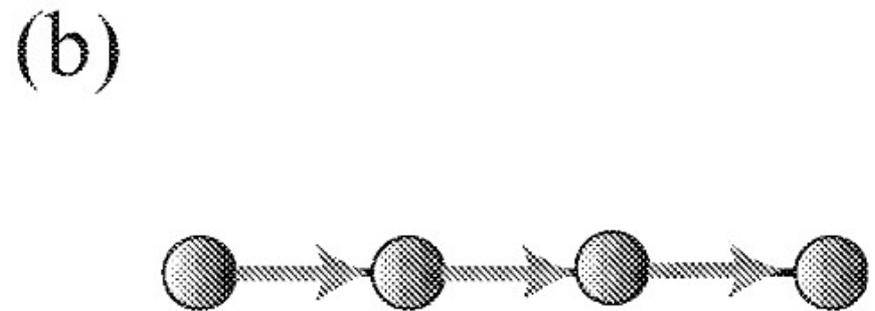
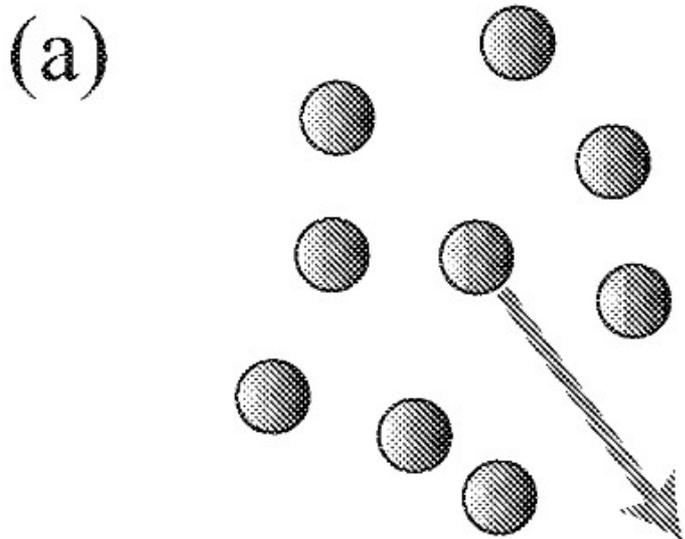
Simone Barbarino, Davide Rossini, Leonardo Mazza, SNS-NEST, Pisa



Rosario Fazio
NEST-SNS
ICTP, Trieste

Physics in 1D

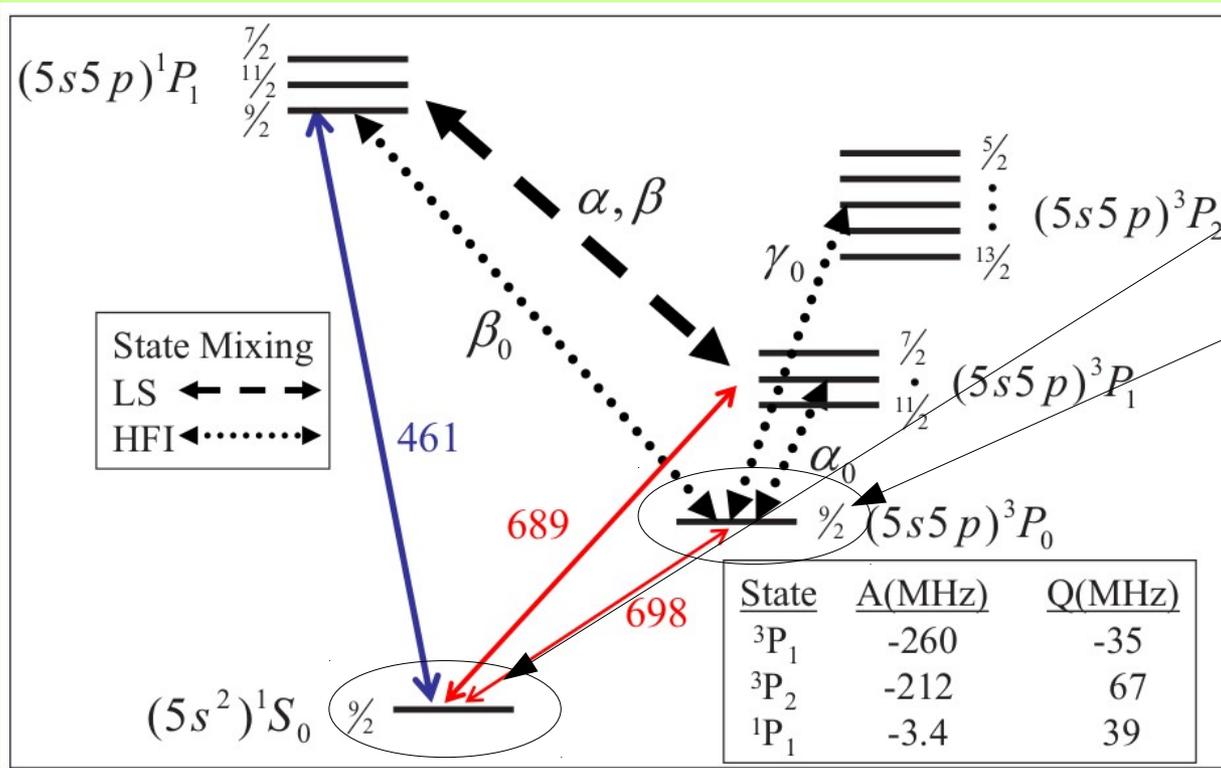
- **Interactions** are intrinsically strong \rightarrow *collective* excitations
- **Statistics**: bosons and fermions behave in a similar way



Peculiar phenomena:

- Long-range order \rightarrow *quasi*-long-range order (power laws)
- Solvability/integrability
- Lack of thermalization
- “Fermionization” of bosons / “bosonization” of fermions
- Anderson/many-body localization
- Spin-charge separation
- ...

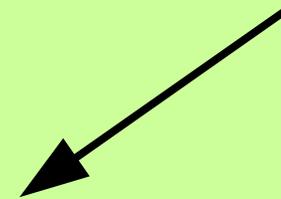
Earth-alkaline(-like) Fermi atoms



^{87}Sr , from Boyd *et al.*, PRA **76**, 022510 (2007)

1) 1S_0 and 3P_0 coupled through an ultranarrow forbidden transition

2) Purely *nuclear spin* → s-wave scattering lengths don't depend on it



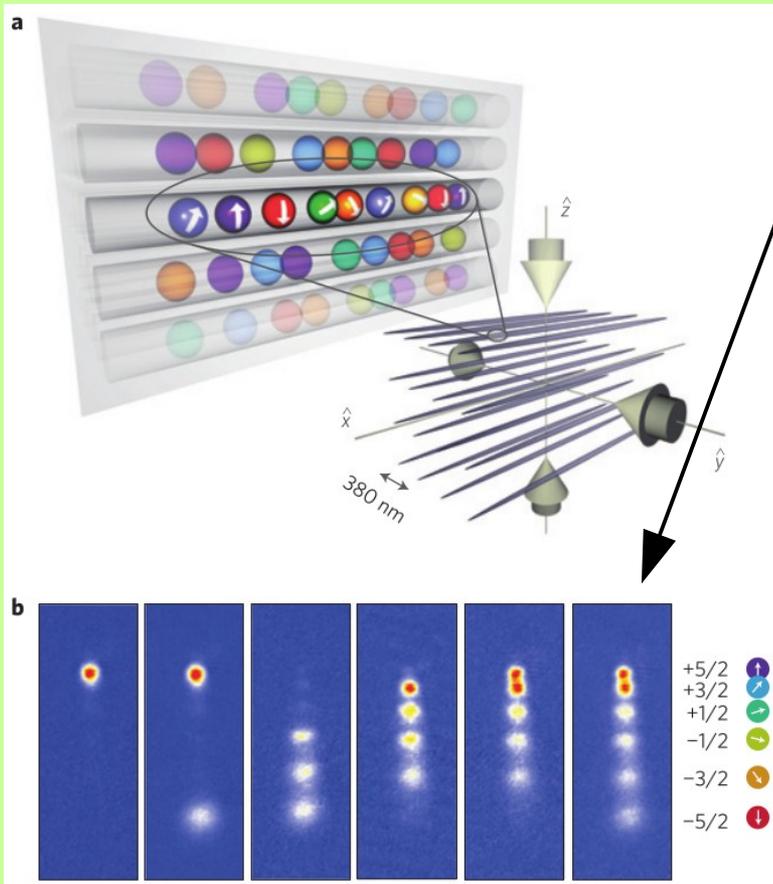
Promising platform to study physics of multi-species gases!

A first experiment in the continuum

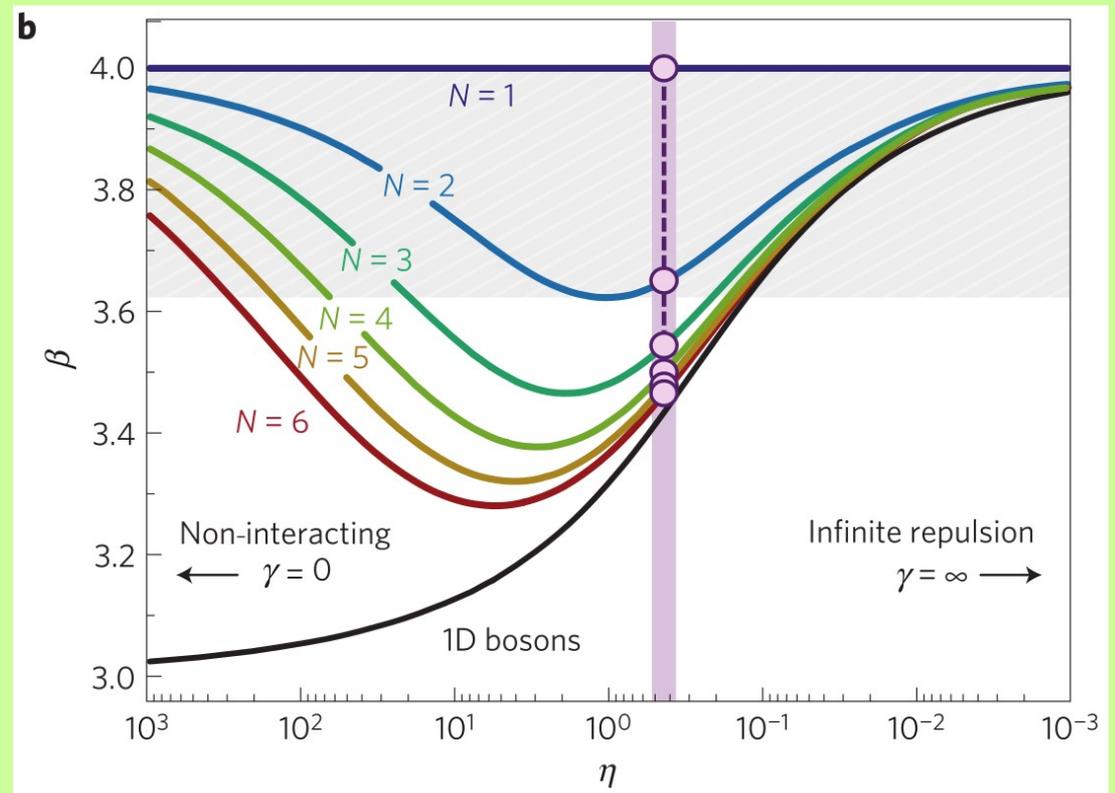
Pagano *et al.*, Nat. Phys. 10, 198 (2014)

^{173}Yb : $l = 5/2$

- The population of the 3P_0 orbital can be neglected
- Species (nuclear-spin m_z) can be populated at will



“Bosonization” of fermions!



Addition of the optical lattice

Gorshkov *et al.*, Nat. Phys. **6**, 289 (2010)

$$H = - \sum_{\alpha=g,e} t_{\alpha} \sum_j \sum_{m=-I}^I (c_{\alpha,j,m}^{\dagger} c_{\alpha,j+1,m} + \text{H. c.}) + \sum_{\alpha=g,e} \frac{U_{\alpha\alpha}}{2} \sum_j n_{\alpha,j} (n_{\alpha,j} - 1) +$$

$$+ V \sum_j n_{g,j} n_{e,j} + V_{ex} \sum_{j,m,m'} c_{g,j,m}^{\dagger} c_{e,j,m'}^{\dagger} c_{g,j,m'} c_{e,j,m}$$

hopping amplitudes direct interaction exchange interaction on-site interaction

Experimental realization (Taie *et al.*, Nat. Phys. **8**, 825 (2012), ¹⁷³Yb):

SU(N) Hubbard model

$$H_0 = -t \sum_j \sum_{m=-I}^I (c_{j,m}^{\dagger} c_{j+1,m} + \text{H. c.}) + U \sum_j \sum_{m < m'} n_{j,m} n_{j,m'}$$

- Assaraf *et al.*, PRB **60**, 2299 (1999)
- Szirmai & Sólyom, PRB **71**, 205108 (2005)
- Buchta *et al.*, PRB **75**, 155108 (2007)
- Manmana *et al.*, PRA **84**, 043601 (2011)

...

Observation of SU(6) Mott insulator!

Synthetic dimension

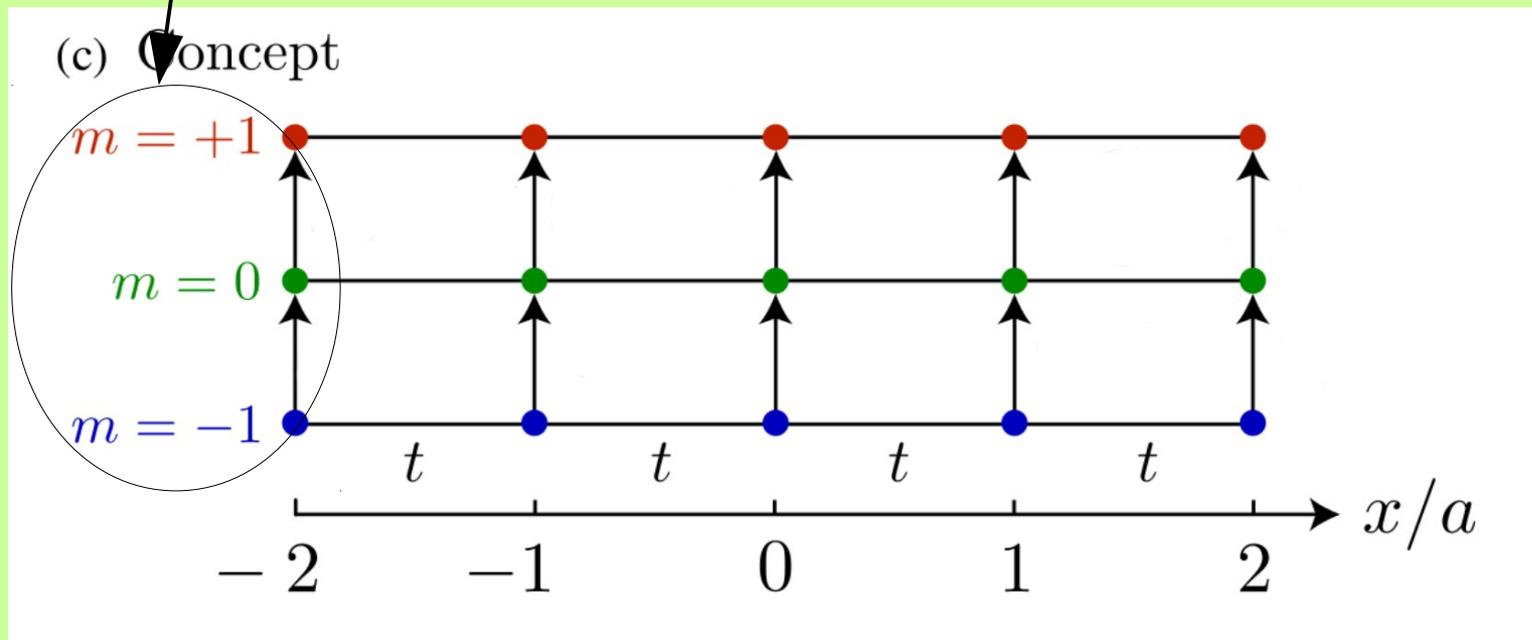
Boada *et al.*, PRL **108**, 133001 (2012)

Idea: use an internal degree of freedom of the atom (in our case, nuclear spin) as an additional dimension

Raman coupling of nuclear-spin levels \rightarrow tunneling in the synthetic dimension

Quantum ladder!

quasi-1D system \rightarrow exploration of physics in 2D



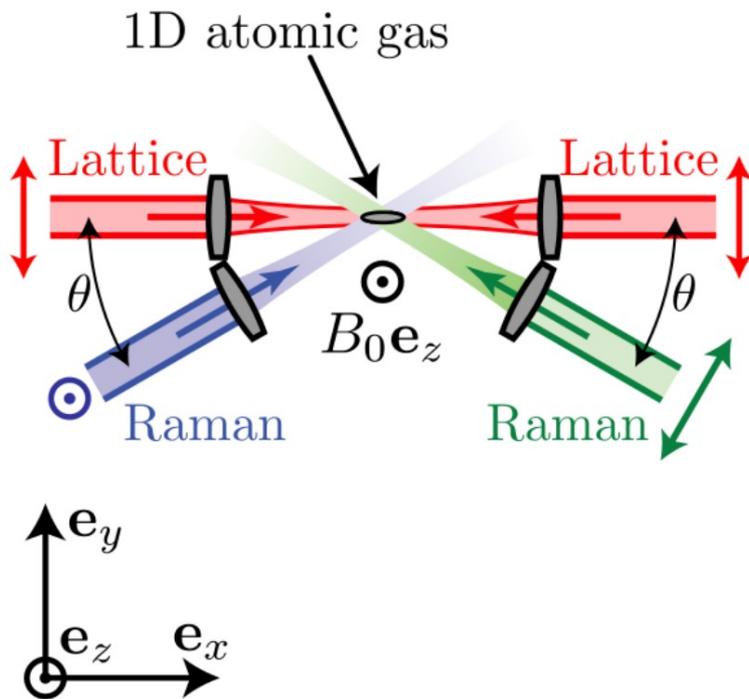
Synthetic gauge field (I)

Celi *et al.*, PRL **112**, 043001 (2014)

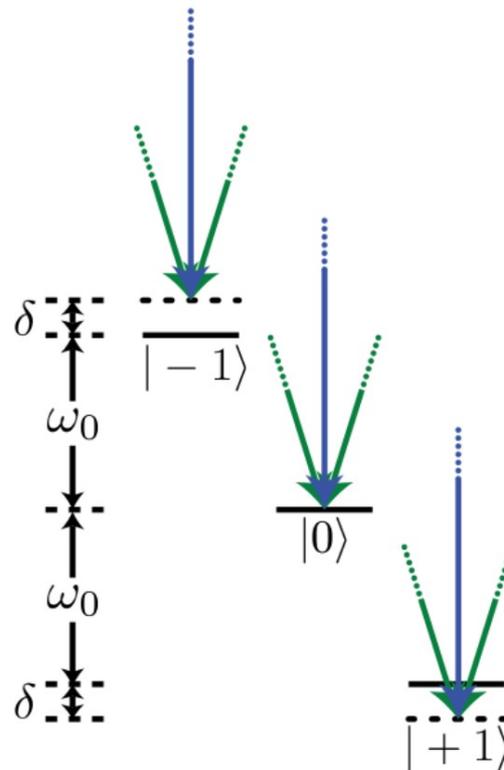
Question: is it possible to engineer a **magnetic-like flux** in the synthetic lattice?

Answer: yes! Raman-assisted coupling in the synthetic dimension

(a) Layout



(b) Level diagram

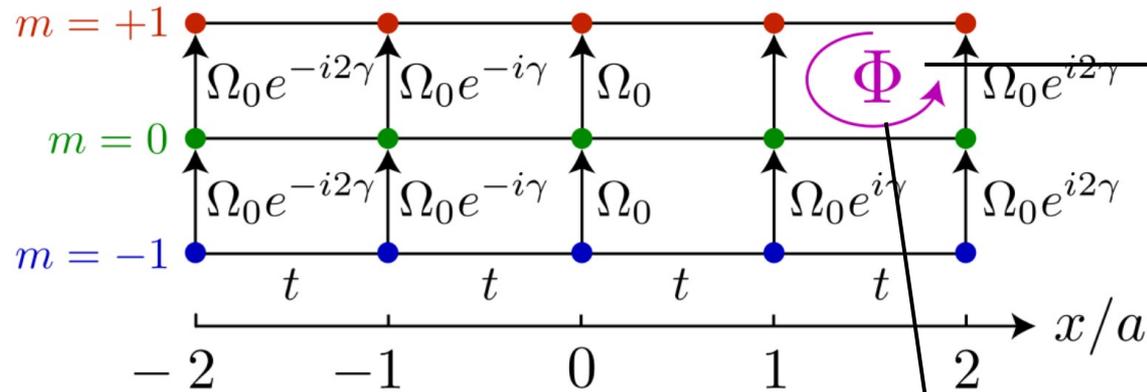


(For other ways of implementing synthetic gauge fields, see, e.g., Dalibard *et al.*, RMP **83**, 1523 (2011))

Synthetic gauge field (II)

Celi *et al.*, PRL **112**, 043001 (2014)

(c) Concept



Magnetic-like flux

$$\Phi = \gamma$$

$$H = H_0 + H_1 + H_2$$

(a **cyclic** coupling between $m = \pm I$ can in principle also be implemented)

$$H_0 = -t \sum_j \sum_{m=-I}^I (c_{j,m}^\dagger c_{j+1,m} + \text{H. c.}) + U \sum_j \sum_{m < m'} n_{j,m} n_{j,m'}$$

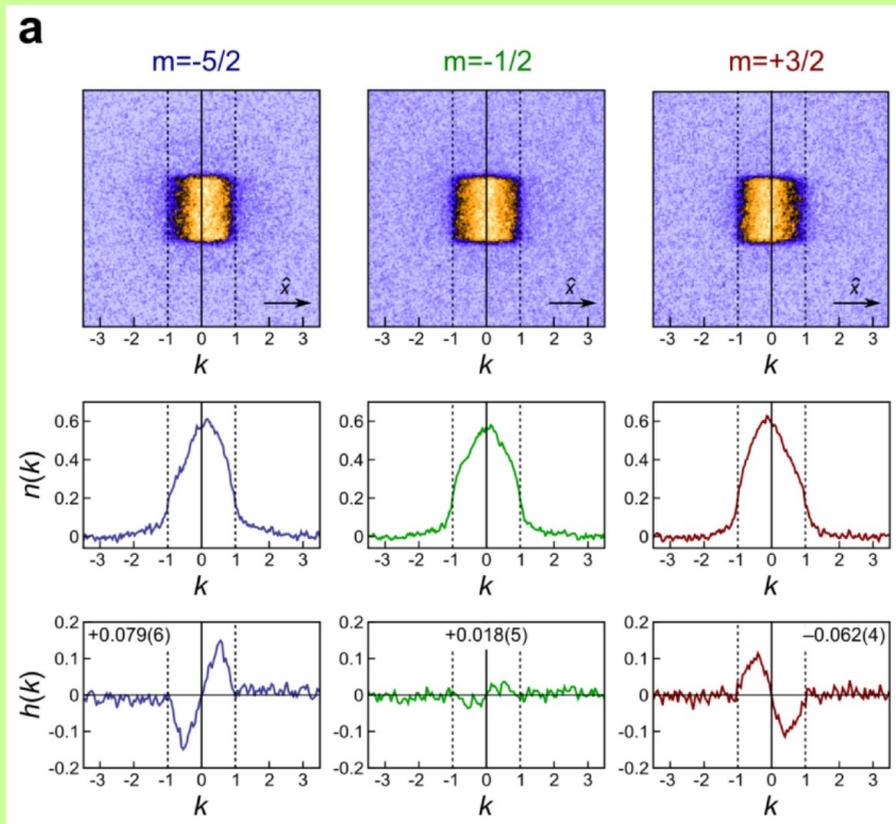
$$H_1 = \Omega \sum_j \sum_{m=-I}^{I-1} (e^{-i\gamma j} c_{j,m}^\dagger c_{j,m+1} + \text{H. c.})$$

$$H_2 = \Omega' \sum_j (e^{-i\gamma j} c_{j,I}^\dagger c_{j,-I} + \text{H. c.})$$

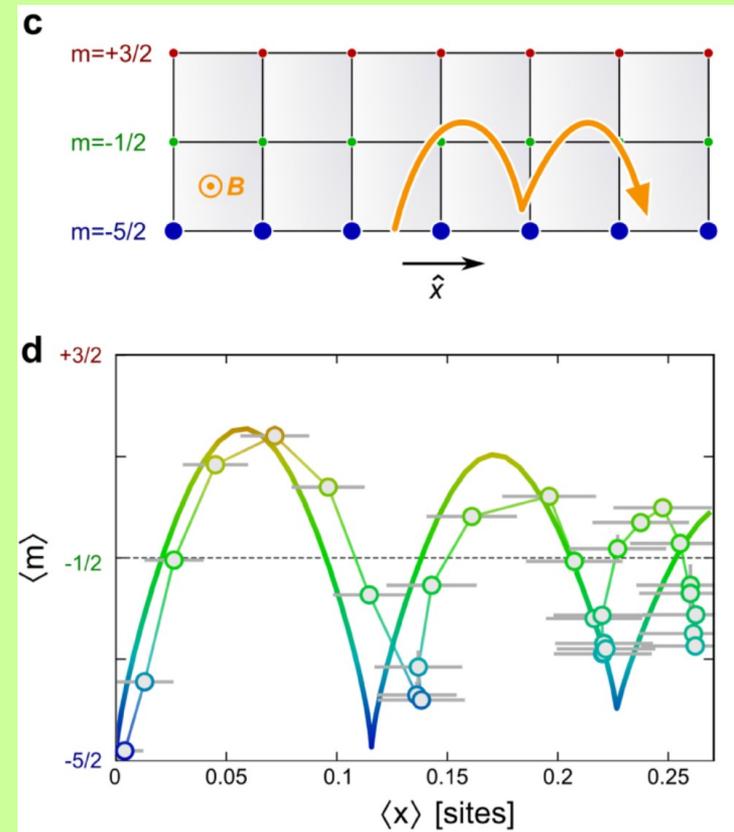
related to the parameters of Raman lasers

A look to experiments

Mancini *et al.*, Science 349, 1510 (2015)



Momentum distributions



Skipping orbits

Evidence for chiral edge modes!

Quantum Hall effect

System: 2D electrons in a magnetic field

Observation: at strong magnetic fields, Hall conductance is quantized

$$\sigma_H = \frac{p}{q} \frac{e^2}{h}$$

$q = 1$: integer QHE; $q > 1$: fractional QHE

Key quantity: *filling fraction*

number of electrons

$$\nu = \frac{N}{N_\Phi} = \frac{p}{q}$$

number of elementary fluxes

Microscopically: the bulk is an *insulator*, and all the conductance is concentrated on the edges

Edge modes!

Integer QHE: *non-interacting physics*

Fractional QHE: Coulomb interaction

Liquid: Laughlin wave functions, at $\nu = \frac{1}{q \text{ odd}}$

Integer QHE with cold atoms

$$H_0 = -t \sum_j \sum_{m=-I}^I (c_{j,m}^\dagger c_{j+1,m} + \text{H. c.}) + U \sum_j \sum_{m < m'} n_{j,m} n_{j,m'}$$

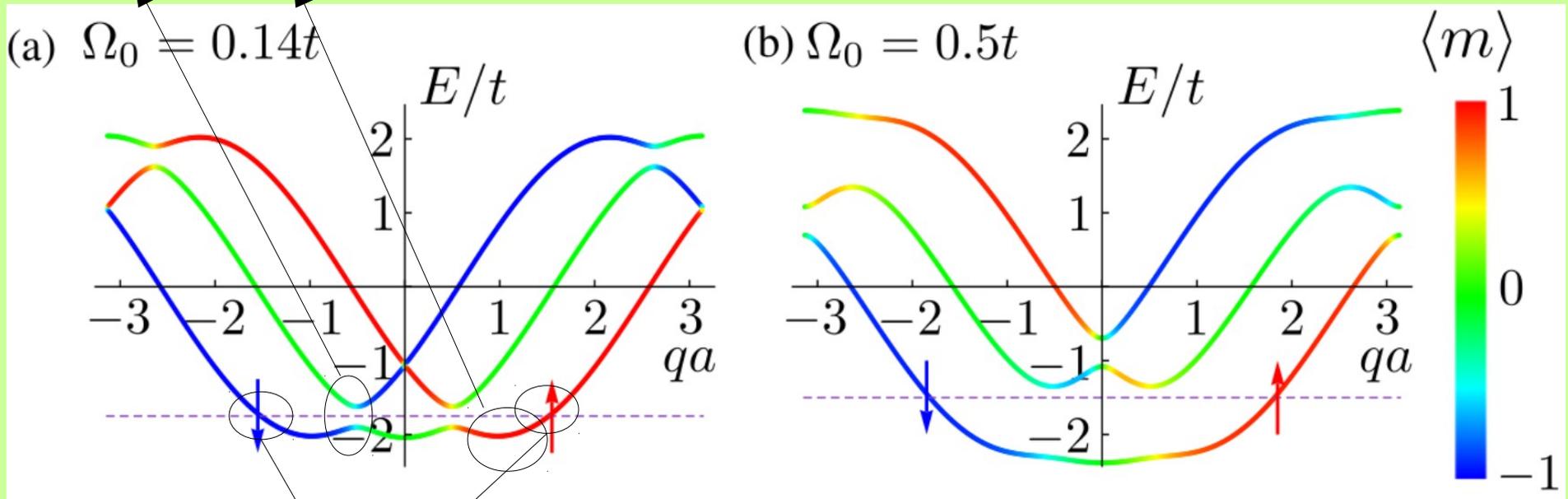
$$H_1 = \Omega \sum_j \sum_{m=-I}^{I-1} (e^{-i\gamma j} c_{j,m}^\dagger c_{j,m+1} + \text{H.c.})$$

$$H_2 = \Omega' \sum_j (e^{-i\gamma j} c_{j,I}^\dagger c_{j,-I} + \text{H.c.})$$

Filling fraction:

$$\nu = \frac{N}{L(2I+1) \frac{\gamma}{2\pi}} \in \mathbb{N}$$

Single-particle spectrum:



From Celi *et al.*, PRL **112**, 043001 (2014)

Helical modes: gapless low-energy excitations with definite momentum and spin!

→ Edge modes of QHE

Questions

#0 To what extent the system is a quantum simulator of QHE?



#1 What happens when interactions are added?

#2 What happens when the cyclic coupling in the synthetic direction is added?

#3 Which observables can characterize the properties of the system?

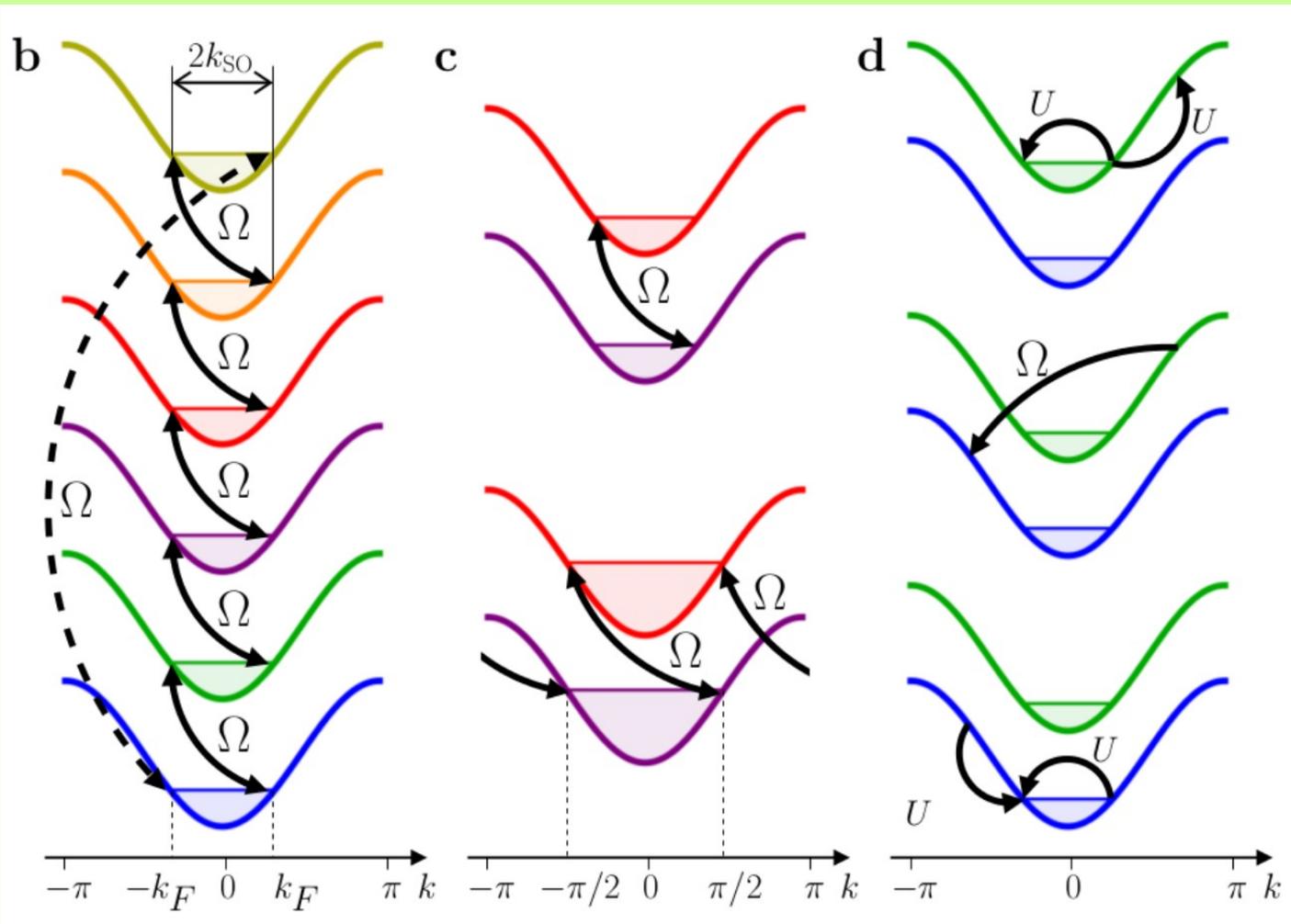
Cyclic coupling

$$H_0 = -t \sum_j \sum_{m=-I}^I (c_{j,m}^\dagger c_{j+1,m} + \text{H. c.}) + U \sum_j \sum_{m < m'} n_{j,m} n_{j,m'}$$

$$H_1 = \Omega \sum_j \sum_{m=-I}^{I-1} (e^{-i\gamma j} c_{j,m}^\dagger c_{j,m+1} + \text{H.c.})$$

$$H_2 = \Omega' \sum_j (e^{-i\gamma j} c_{j,I}^\dagger c_{j,-I} + \text{H.c.})$$

Each term in H_1 opens a gap $\sim \Omega$ in the single particle spectrum \rightarrow in the non-interacting case, the system becomes fully gapped at fillings $\nu \in \mathbb{N} \rightarrow N(\nu, \gamma)$

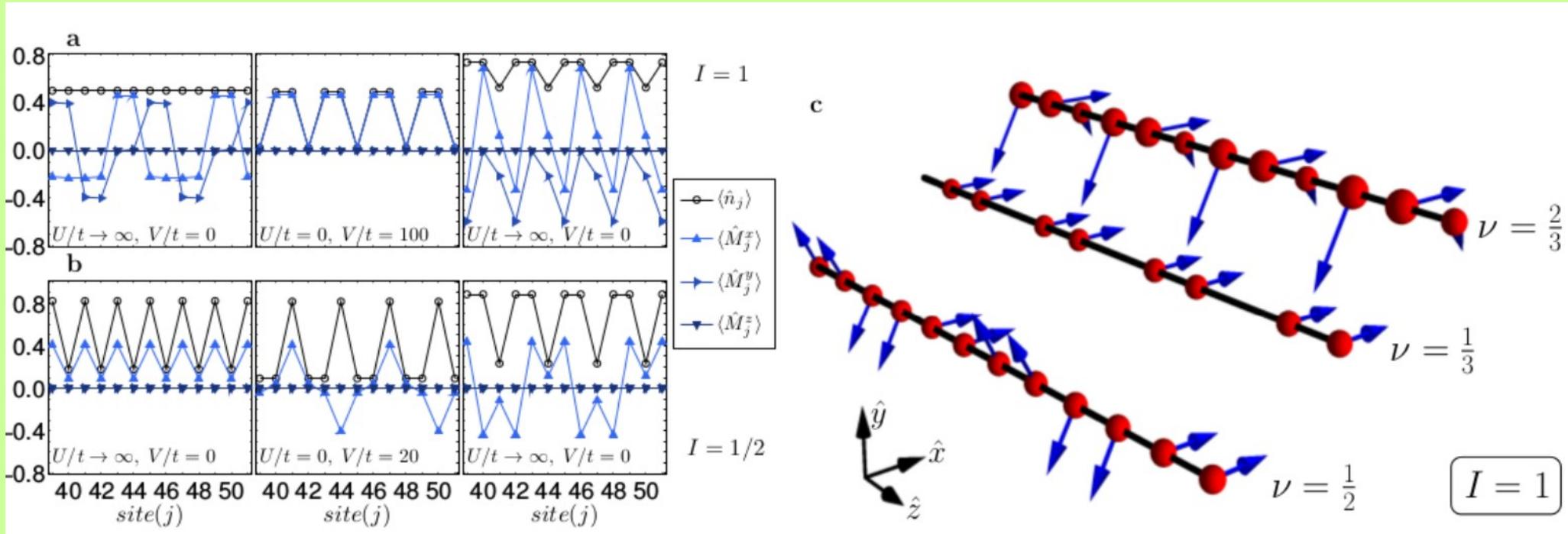


When $\Omega, U \neq 0$, a many-body gap opens at fillings $\nu \in \mathbb{Q}$

Fractional QHE!

Magnetic crystals #1

Computational method: **DMRG** (White, PRL **69**, 2863 (1992))



Crystalline order!

“Problem”: decreasing ν ,
longer-ranged interactions are needed in order to stabilize the crystal \longrightarrow **Dipoles?**

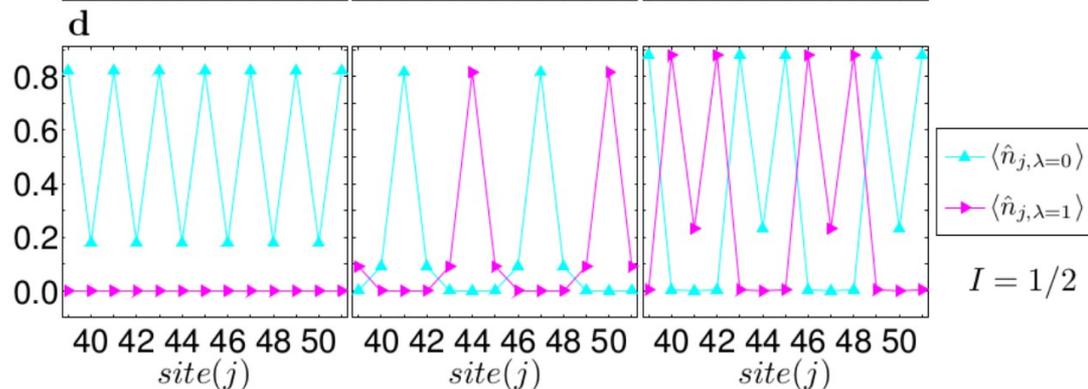
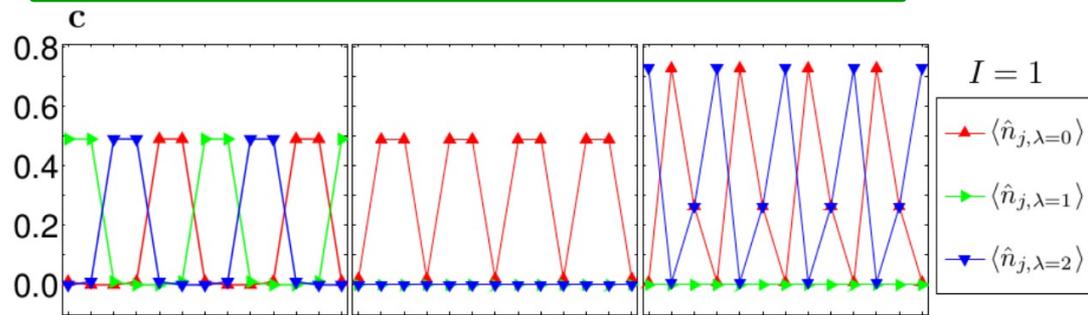
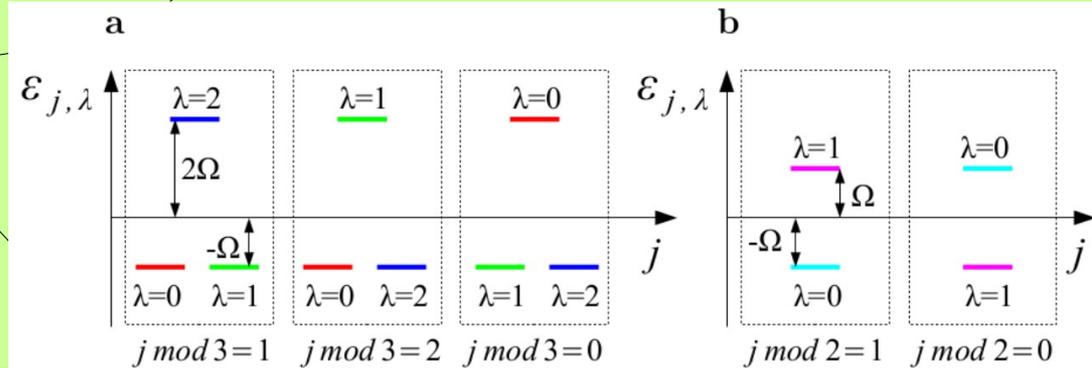
Magnetic crystals #2

$$H_0 = -t \sum_j \sum_{m=-I}^I (c_{j,m}^\dagger c_{j+1,m} + \text{H. c.}) + U \sum_j \sum_{m < m'} n_{j,m} n_{j,m'}$$

$$H_1 = \Omega \sum_j \sum_{m=-I}^{I-1} (e^{-i\gamma j} c_{j,m}^\dagger c_{j,m+1} + \text{H.c.})$$

$$H_2 = \Omega' \sum_j (e^{-i\gamma j} c_{j,I}^\dagger c_{j,-I} + \text{H.c.})$$

True structure of the crystals!



Diagonalization (Fourier)

$$\begin{cases} d_{j,\lambda} = \frac{1}{\sqrt{2I+1}} \sum_m e^{i\frac{2\pi\lambda}{2I+1}m} c_{j,m} \\ \epsilon_{j,\lambda} = 2\Omega \cos\left(\frac{2\pi\lambda}{2I+1} + \gamma j\right) \end{cases}$$

Helical liquids

$$H_0 = -t \sum_j \sum_{m=-I}^I (c_{j,m}^\dagger c_{j+1,m} + \text{H. c.}) + U \sum_j \sum_{m < m'} n_{j,m} n_{j,m'}$$

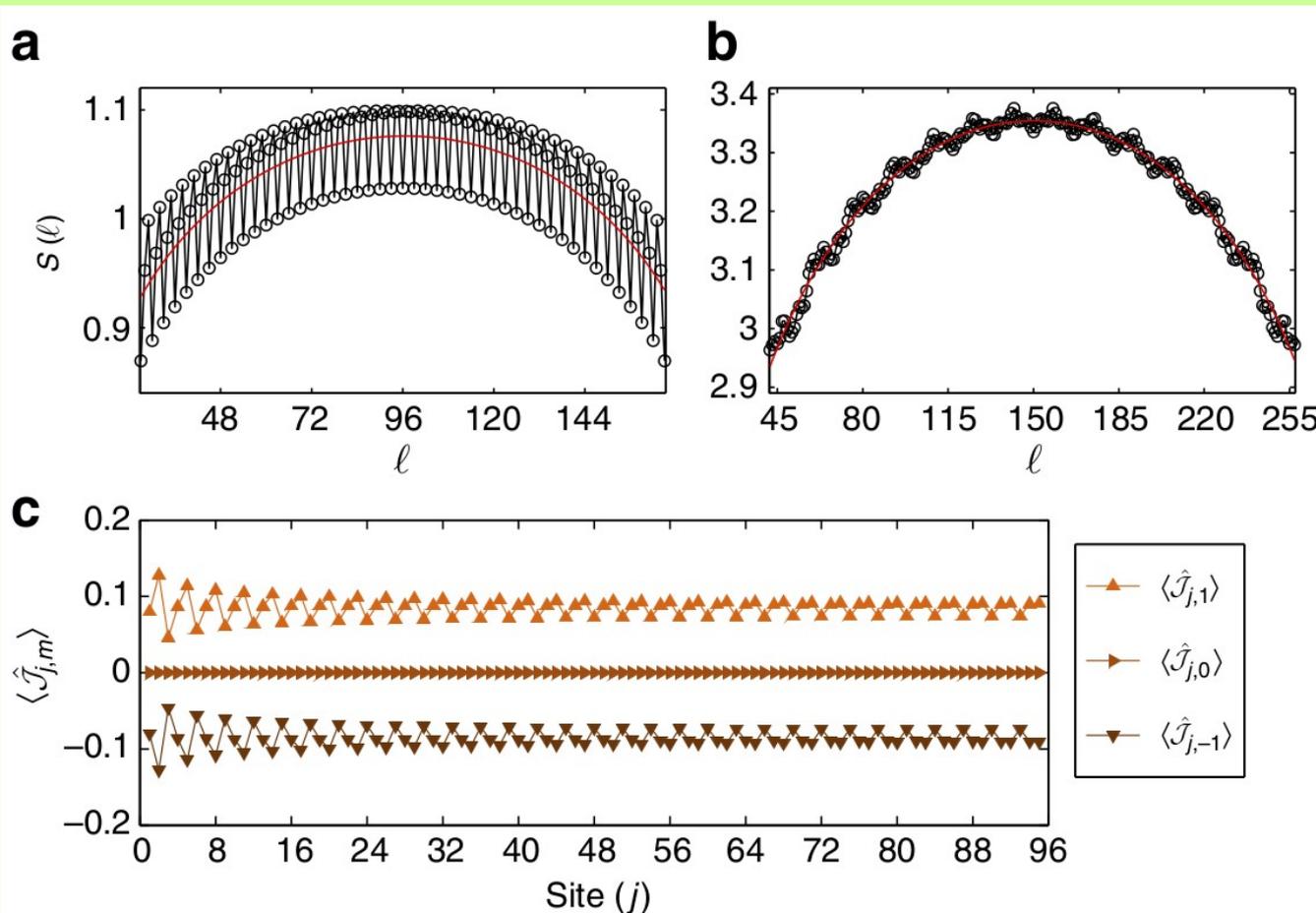
$$H_1 = \Omega \sum_j \sum_{m=-I}^{I-1} (e^{-i\gamma j} c_{j,m}^\dagger c_{j,m+1} + \text{H.c.})$$

$$\cancel{H_2 = \Omega' \sum_j (e^{-i\gamma j} c_{j,I}^\dagger c_{j,-I} + \text{H.c.})}$$

Helical liquids should survive!

Spin-resolved currents + entanglement entropies

central charge



Filling, Raman coupling and interactions are all necessary in order to stabilize the liquids!

Helicity and interactions

Questions:

Is helicity affected by **interactions**?

Can we find experimentally detectable quantities that witness chirality?

Momentum distribution functions

$$n_{p,m}, \quad n_p = \sum_m n_{p,m}$$

Mean current

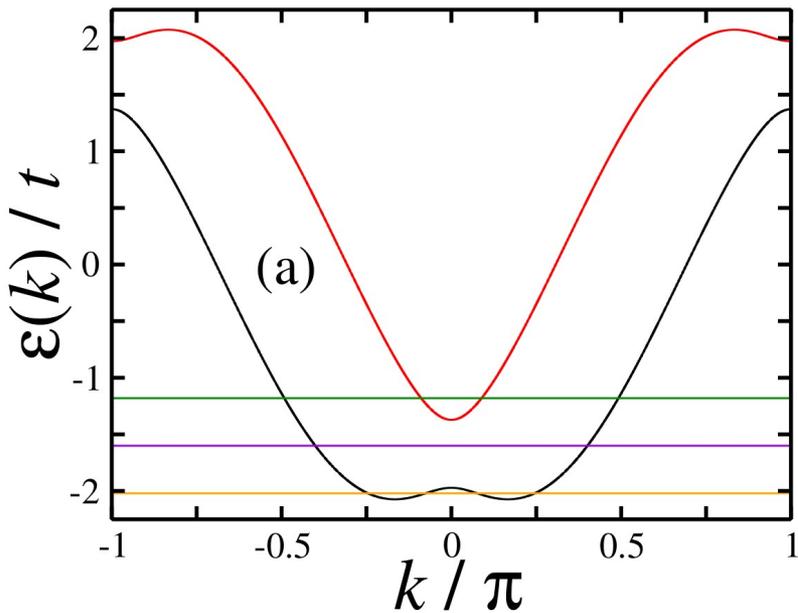
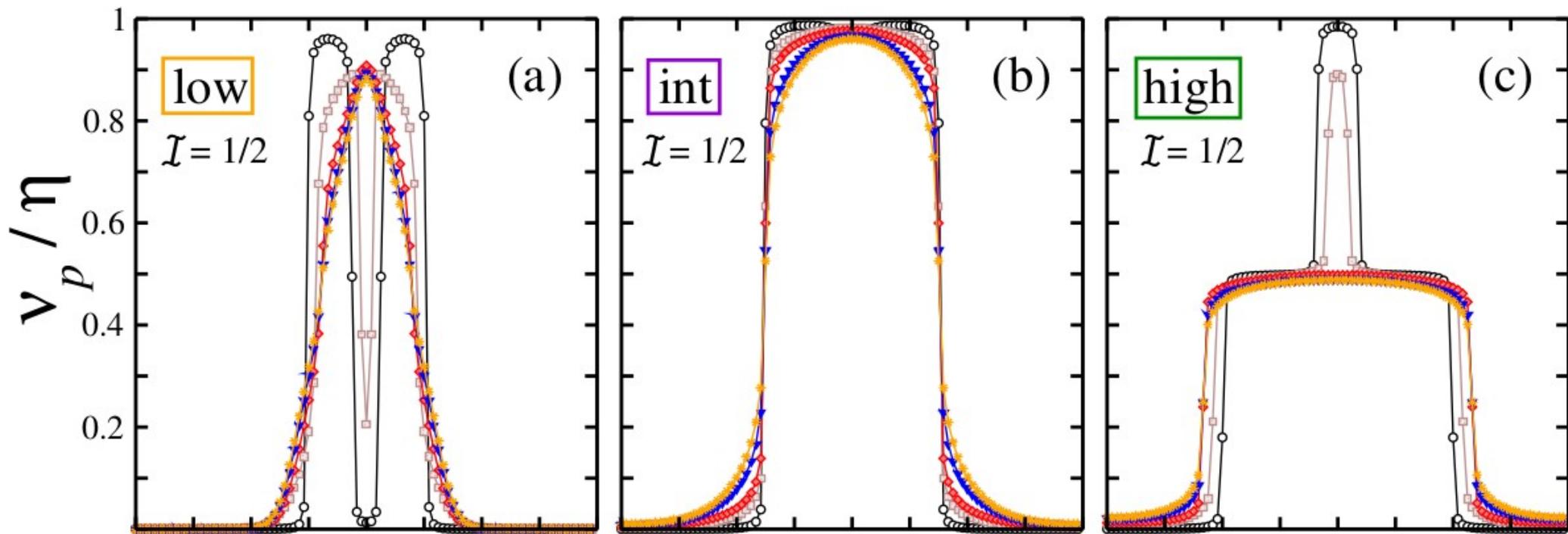
$$Q_m = \frac{1}{L} \sum_j C_{j,m} = \frac{-2t}{L} \sum_{p>0} (n_{p,m} - n_{-p,m})$$

Asymmetry of the momentum distributions

$$J_m = - \sum_{p>0} (n_{p,m} - n_{-p,m})$$

**Key observables:
momentum
distributions!**

Momentum distributions #1

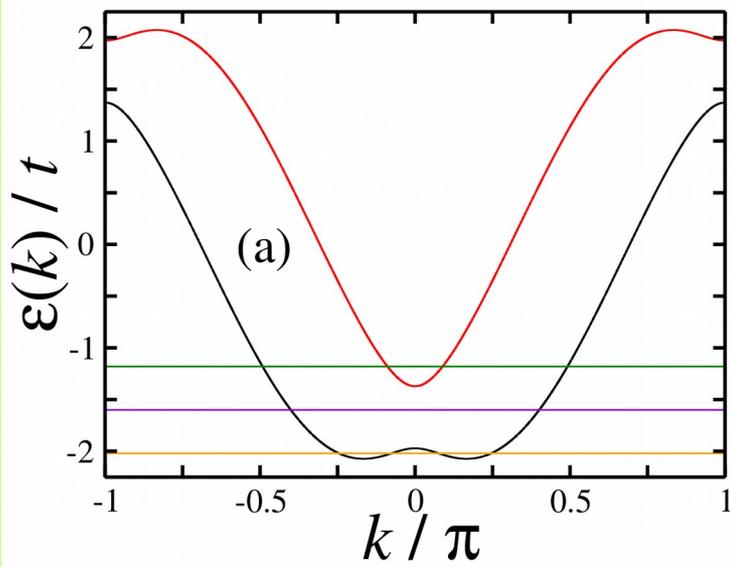
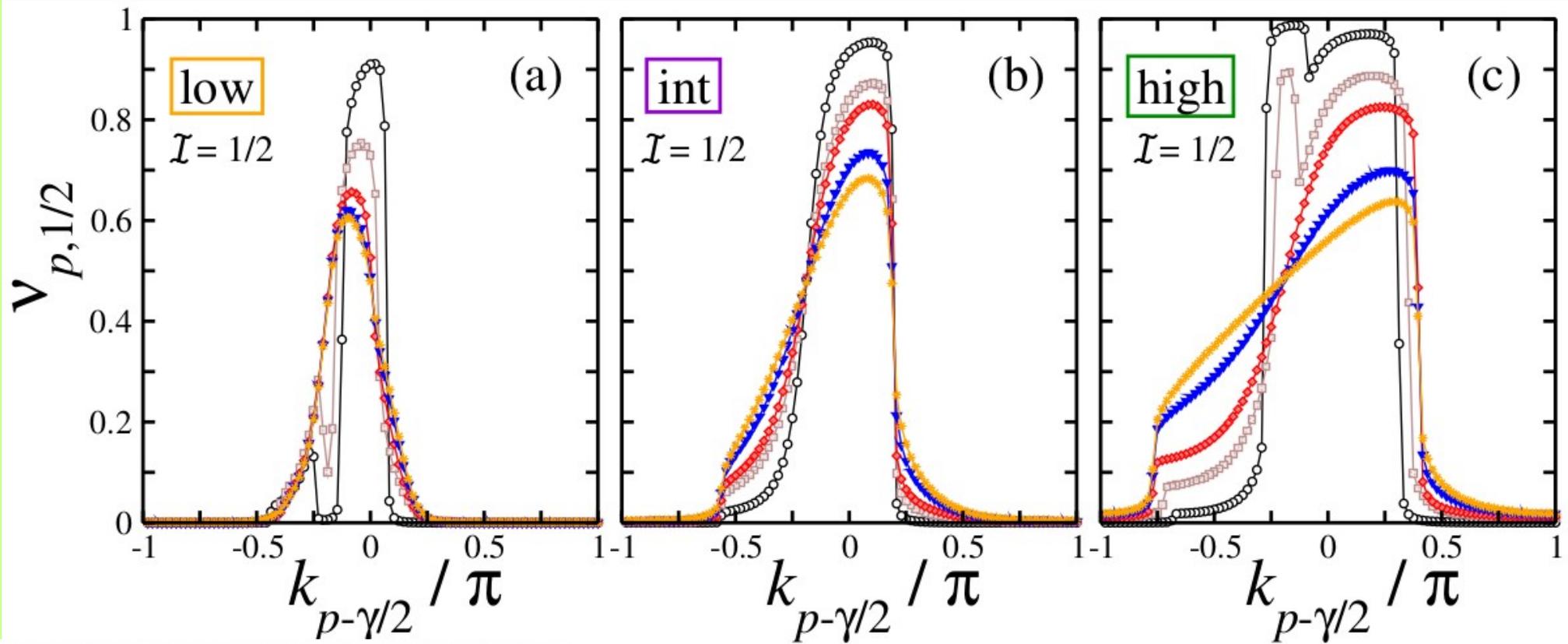


The total momentum distributions show that increasing U results in an **effective increasing of Ω**



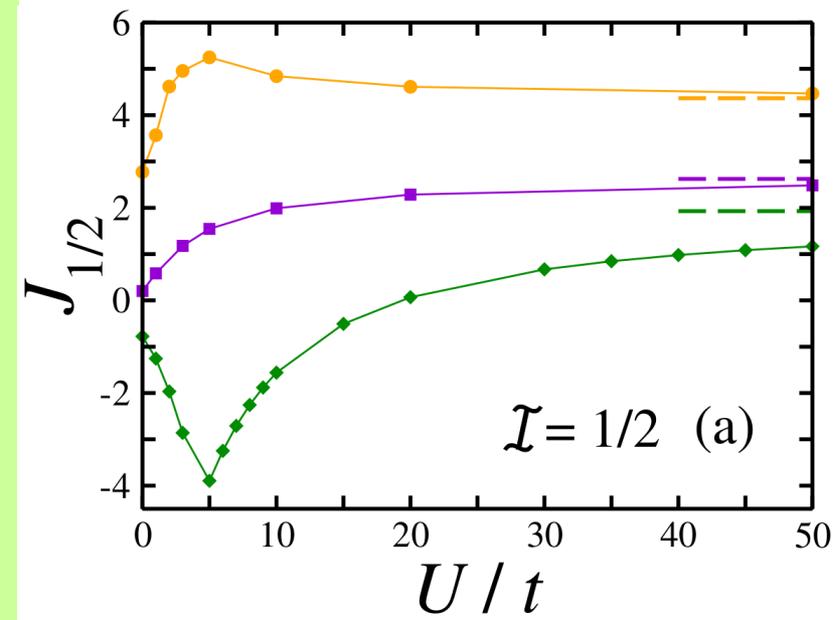
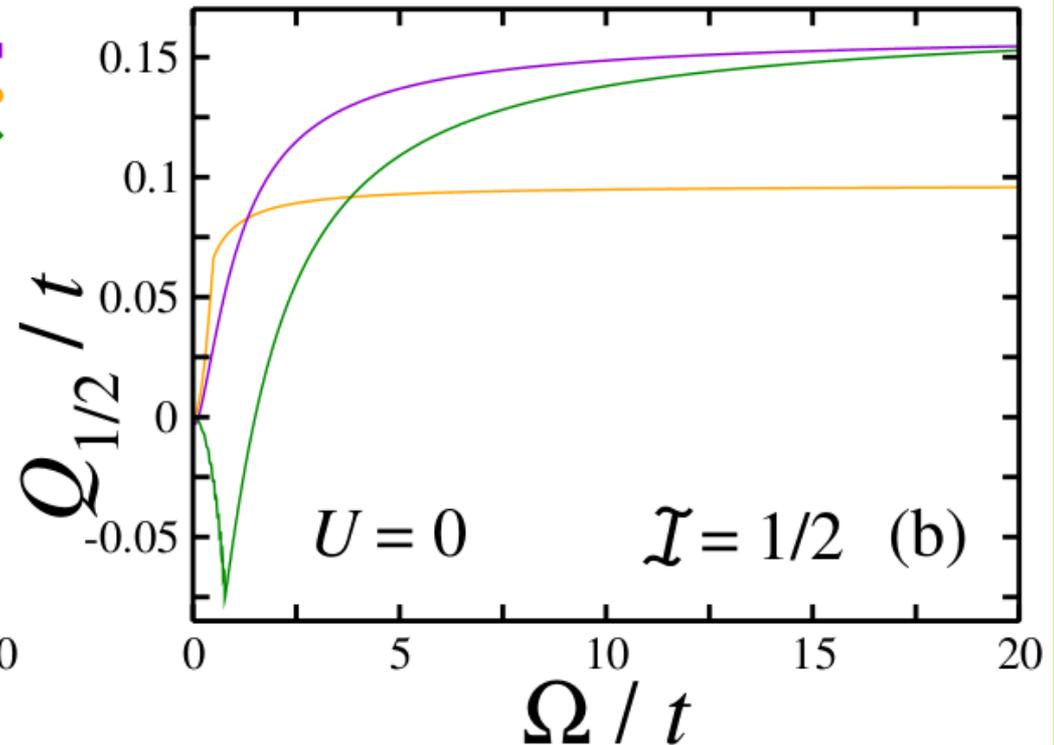
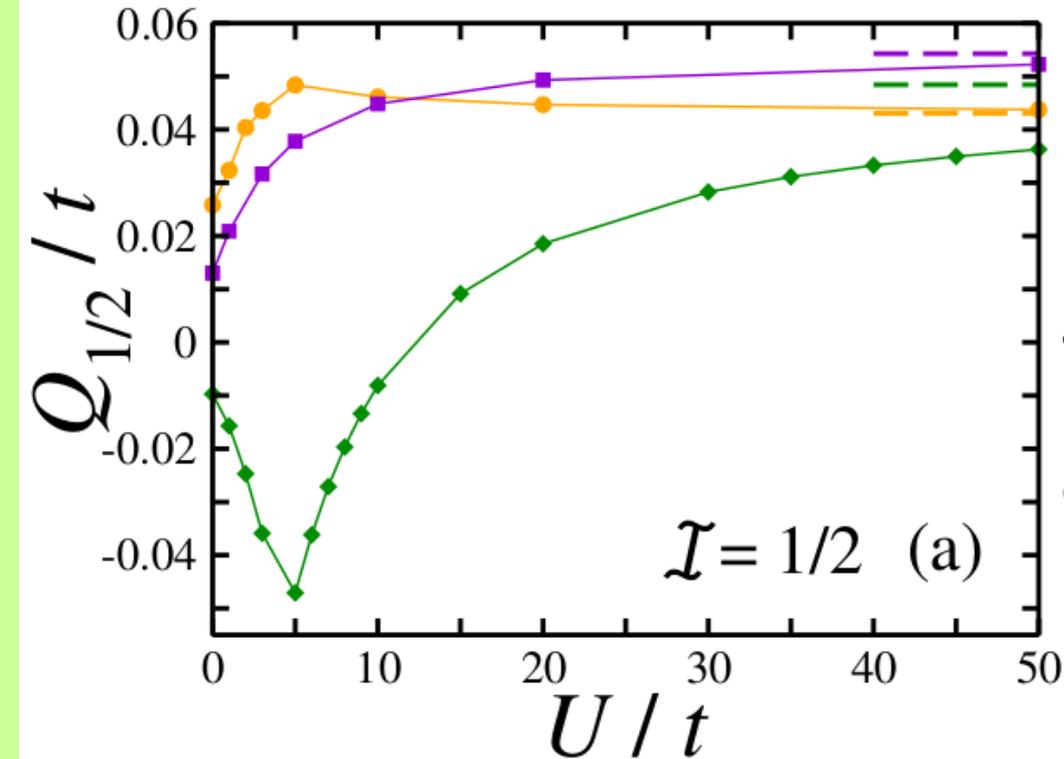
Theoretical prediction from Braunecker *et al.*, PRB **82**, 045127 (2010)

Momentum distributions #2



The shape of the spin-resolved momentum distributions changes increasing $U \rightarrow$ **asymmetry wrt $k = 0$ is enhanced**

Asymmetry and mean current



- Mean currents display non-trivial behaviors very similar each other
- The behaviors are explained by the effective enhancement of Ω

Conclusions and outlook

- Earth-alkaline-like atoms offer a promising platform for the study of the QHE in cold atoms
- We showed that some features of QHE emerge when a synthetic gauge field is implemented in the lattice gas
- We show how helicity witnesses are affected by interactions

Some other interesting directions:

- In which regimes is the system QHE-like? → Phase diagram
- What is the role of the form of the interaction for the QHE physics?
- Is the system topological?
- Does the system possess fractional/anionic excitations?
- ...

Thank you for your attention!